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Pierre Valla, Sally Lowick, Frédéric Herman, Jean-Daniel Champagnac, Philippe Steer, et al.. Exploring IRSL 50 fading variability in bedrock feldspars and implications for OSL thermochronometry. *Quaternary Geochronology*, 2016, 36, pp.55-66. 10.1016/j.quageo.2016.08.004 . insu-01355727

HAL Id: insu-01355727

<https://hal-insu.archives-ouvertes.fr/insu-01355727>

Submitted on 24 Aug 2016

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Accepted Manuscript

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PII: S1871-1014(16)30104-2

DOI: [10.1016/j.quageo.2016.08.004](https://doi.org/10.1016/j.quageo.2016.08.004)

Reference: QUAGEO 784

To appear in: *Quaternary Geochronology*

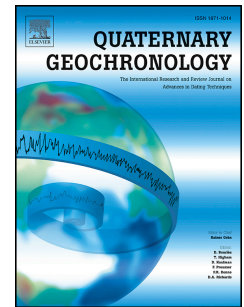
Received Date: 21 December 2015

Revised Date: 22 July 2016

Accepted Date: 8 August 2016

Please cite this article as: Valla, P.G., Lowick, S.E., Herman, F., Champagnac, J.-D., Steer, P., Guralnik, B., Exploring IRSL₅₀ fading variability in bedrock feldspars and implications for OSL thermochronometry, *Quaternary Geochronology* (2016), doi: 10.1016/j.quageo.2016.08.004.

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Exploring IRSL₅₀ fading variability in bedrock feldspars and implications for OSL thermochronometry

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Abstract

Optically Stimulated Luminescence (OSL) is a well-established Quaternary dating method, which has recently been adapted to application in low-temperature thermochronometry. The Infra-Red Stimulated Luminescence (IRSL) of feldspar, which so far is the most promising target signal in thermochronometry, is unfortunately prone to anomalous fading. The fading of feldspar IRSL is at times not only challenging to measure, but also laborious to incorporate within luminescence growth models. Quantification of IRSL fading is therefore a crucial step in OSL thermochronometry, raising questions regarding (i) reproducibility and reliability of laboratory measurements of fading, as well as (ii) the applicability of existing fading models to quantitatively predict the level of IRSL field saturation in nature. Here we investigate the natural luminescence signal and anomalous fading of IRSL measured at 50 °C (IRSL₅₀) in 32 bedrock samples collected from a variety of lithologies and exhumation settings (Alaska and Norway). We report a large span of IRSL₅₀ fading rates between samples ($g_{2\text{days}}$ ranging from ~0.5 to ~45 %/decade), which further demonstrates (i) a good reproducibility between two common fading measurement protocols, and (ii) the ability of tunnelling models to predict the level of feldspar IRSL₅₀ field saturation in nature. We observe higher IRSL₅₀ fading in feldspar with increasing Ca content, although other factors cannot be dismissed at present. Finally, our dataset confirms that the applicability of feldspar IRSL₅₀ in OSL thermochronometry is limited to rapidly-exhuming settings or warm subsurface environments.

Keywords: luminescence, feldspar IRSL, anomalous fading, OSL thermochronometry

1. Introduction

1.1 OSL thermochronometry

Understanding the processes operating at the Earth's surface requires quantitative methods to measure erosion and sedimentation rates. In that context, low-temperature thermochronometric methods (see Reiners and Brandon, 2006 for a review) have been increasingly used over the past decades to quantify bedrock cooling histories within the top ~2-8 km of the upper crust, and subsequently to interpret them into exhumation histories (e.g. Braun et al., 2006; 2012). However, linking long-term exhumation rates (10^5 - 10^8 year timescales) to short-term erosion measurements (10^0 - 10^3 year timescales) remains challenging, because different spatial scales and processes are involved. An ongoing debate (Herman and Champagnac, 2016; Willenbring and Jerolmack, 2016), regarding the impact of Quaternary climate and glacial/interglacial oscillations on global-scale erosion rates, highlights the current shortage of quantitative estimates of bedrock erosion and relief evolution over Quaternary to sub-Quaternary timescales (e.g. Herman et al., 2013; Champagnac et al., 2014). Recent methodological developments of low-temperature thermochronometers, such as $^4\text{He}/^3\text{He}$ (Shuster and Farley, 2005) and Optically Stimulated Luminescence (OSL) (Guralnik et al., 2015a; King et al., 2016), offer an unprecedented resolution to constrain rock cooling histories within the upper ~1-3 km of the Earth's crust, and thus to potentially reconstruct erosion histories and/or topographic relief changes over sub-Quaternary timescales (e.g. Shuster et al., 2005; 2011; Valla et al., 2011a; King et al., 2016).

OSL dating (e.g. Aitken, 1998) has recently been adapted for low-temperature thermochronometry by considering the competing effects of filling and emptying of naturally-occurring electron traps (linked to the OSL signal) in crystal lattices due to environmental radiation and temperature, respectively (Herman et al., 2010; Li and Li, 2012; Guralnik et al., 2013; King et al., 2016). While other "trapped charge dating" methods (e.g. Grün, 2001) have been proposed and tested, including thermoluminescence (TL; e.g. Prokein and Wagner, 1994; Tang and Li, 2015) and electron-spin resonance (ESR; Koshchug and Solovyov, 1998; Grün et al., 1999), the advantage of OSL techniques in thermochronometry lies in their rapid methodological development and expanding use in Quaternary science (see reviews by Wintle, 2008 and Rhodes, 2011). One major intrinsic limitation in trapped charge dating, common to OSL, ESR and TL alike, is signal saturation (i.e. when all available electron traps become

filled). If the target luminescence signal for dating is close to its saturation limit, only minimum ages may be provided. Consequently, when luminescence is used to decipher rock cooling rates, near-saturation signals can only characterise maximum (upper-bound) rates of cooling and exhumation (e.g. Guralnik et al., 2013). Fortunately, the first OSL-thermochronometry applications (Herman et al., 2010; Wu et al., 2012; De Sarkar et al., 2013; Wu et al., 2015) addressed settings of extremely rapid exhumation (where signal saturation should not be a problem), and quantified bedrock cooling histories below ~ 100 °C and over sub-Quaternary timescales. Despite the fact that a standard mineral/signal pair was chosen (quartz “fast OSL component”; Wintle and Murray, 2006), the low sensitivity (Herman et al., 2010; Jeong and Choi, 2012), anomalous characteristics (Li and Li, 2006; Preusser et al., 2009), and the ubiquitous lack of the “fast component” in bedrock quartz of different petrology/mineralogy (Guralnik et al., 2015b), appeared to narrow the use of bedrock quartz for OSL-thermochronometry, and stimulated the continued search for alternative mineral/signal combinations (Guralnik et al., 2015b).

Infra-Red Stimulated Luminescence (IRSL) from feldspar is advantageous over quartz in several ways (e.g. Huntley and Lamothe, 2001), with (i) a higher luminescence sensitivity in both sediment and bedrock samples (e.g. King et al., 2014), (ii) in general, a higher charge storage capacity (i.e. later saturation) offering a wider dating range (e.g. Buylaert et al., 2012), and (iii) a negligible signal contamination from quartz OSL, if present (e.g. Sohbati et al., 2013). Recently, OSL thermochronometry using bedrock Na-feldspar was investigated in a deep continental drillhole (Guralnik et al., 2015a), where IRSL signals measured at 50 °C (hereafter IRSL₅₀) were successfully characterised in the laboratory and translated into environmental palaeotemperatures in the ~ 40 -70 °C range over a ~ 50 ka timescale (Guralnik et al., 2015a). Moreover, the recently-introduced protocol of multiple elevated temperature post-IR (MET-pIRIR) measurement of feldspar (Li and Li, 2011a,b) conceptually offers access to multiple thermochronometric subsystems within a single mineral, as the different subsignals display higher thermal stabilities with increasing stimulation temperatures (Li and Li, 2011b; King et al., 2016).

1.2. Anomalous fading in feldspar IRSL

Although feldspar IRSL appears in many ways preferable to quartz OSL for thermochronometry, one of the drawbacks of feldspar IRSL is its anomalous fading (e.g. Spooner, 1994; Huntley and Lamothe, 2001), an athermal loss of trapped electrons via quantum

mechanical tunnelling to their nearest recombination centres (Visocekas, 2002; Huntley, 2006). Since anomalous fading is ubiquitous in feldspar IRSL₅₀ (Huntley and Lamothe, 2001; Huntley and Lian, 2006), recent efforts have mainly focused on extracting a more stable feldspar signal using pIRIR protocols (e.g. Thomsen et al., 2008; 2011; Li and Li, 2011a; Buylaert et al., 2012). Although the latter protocols exhibit a reduced, perhaps even negligible, anomalous fading, it is generally at the expense of a higher thermal stability of the target electron trap (Li and Li, 2011a,b; Guralnik et al., 2015c; King et al., 2016). This in turn leads to higher residual doses in sediment dating (e.g. Wintle, 2008), and higher closure temperatures in thermochronometry (e.g. Guralnik et al., 2015c), which may once again limit the benefit of non-fading IRSL signals in low-temperature thermochronometry (e.g. Qin et al., 2015). Moreover, the MET-pIRIR protocol often exhibits progressively lower IRSL sensitivity with increasing stimulation temperature (Li and Li, 2011b; King et al., 2016), which can hinder signal extraction for some bedrock samples with low luminescence sensitivity. Despite the presence of anomalous fading in feldspar IRSL₅₀, and in light of recent success in avoiding it via pulsed stimulation at low temperature (Tsukamoto et al., 2006; Jain et al., 2015a), IRSL₅₀ remains highly relevant for OSL thermochronometric investigations.

Over the years, several protocols have been proposed to quantify the anomalous fading of feldspar over laboratory timescales (Huntley and Lamothe, 2001; Auclair et al., 2003) and to correct for fading in the linear dose response range (e.g. Lamothe and Auclair, 1999; Huntley and Lamothe, 2001; Lamothe et al., 2003). More recently, Huntley (2006) formulated a physical description of electron tunnelling that further enabled the application of fading corrections within the entire dose response range (Kars et al., 2008; Li and Li, 2008). In particular, these models allow the prediction of luminescence intensities in “field saturation” (Huntley and Lian, 2006), in which the natural intensity of feldspar IRSL₅₀ becomes insensitive to time, and is governed only by the counterbalancing effects of environmental radiation and anomalous fading (Kars et al., 2008).

In the present study, we investigate anomalous fading of the IRSL₅₀ signal in bedrock feldspar, and highlight its key implications for OSL thermochronometry. As described earlier, OSL thermochronometry exploits the thermal dependence of trapped electrons in minerals during their exhumation towards the Earth’s surface. Given zero fading, a non-saturated natural IRSL₅₀ signal in bedrock feldspar should, in theory, reflect the sample’s cooling history. Anomalous fading of the signal would, however, lower the measured IRSL₅₀ level in a given

bedrock sample, preventing full saturation of the signal and influencing the determination of a cooling rate (if not accounted for). While the incorporation of the effects of IRSL₅₀ fading in OSL-thermochronometric interpretation has been demonstrated as important (Guralnik et al., 2015a; King et al., 2016), there is a lack of consensus at present, as to whether laboratory measurements of IRSL fading are reproducible and/or reliable (see Li et al., 2014 for a review), and thus whether the use of fading models (and their associated age correction schemes) is appropriate. Moreover, the recent notion that all luminescence processes in feldspars are intrinsically related to fading (Jain et al., 2012; 2015b) raises a renewed interest in the fading model of Huntley (2006) and its later derivatives. More specifically, it raises the need for validation studies, where a laboratory measurement of fading rate can be related to its natural fading rate in nature. To date, only a handful of studies (Huntley and Lian, 2006; Kars et al., 2008; Guralnik et al., 2015a) have quantitatively evaluated the effects of natural versus laboratory fading rates. It is therefore necessary to assess (i) whether anomalous fading of the feldspar IRSL₅₀ signal can be accurately measured, and (ii) for a given fading rate, whether the field saturation of the IRSL₅₀ signal can be quantitatively predicted. To this aim, we extracted a wide range of K- and Ca-/Na-feldspars from 32 bedrock samples in different lithologies and exhumation environments, and measured their natural feldspar IRSL₅₀ intensities alongside the anomalous fading of the laboratory-regenerated signals. A large geochemical dataset, including whole-rock major and trace elements as well as selected feldspar mineralogy, enables us to discuss potential causes for the wide range of IRSL₅₀ anomalous fading observed. Finally, we use this dataset to test the theoretical ability of tunnelling models (Huntley 2006; Kars et al., 2008; Guralnik et al., 2015a,c) to predict IRSL₅₀ field saturation levels, and to discuss the implications for OSL thermochronometry.

2. Samples and methods

2.1. Sampling strategy and sample preparation

Bedrock samples from old crystalline or metamorphic basements have been collected in different exhumation settings in Norway and Alaska to explore the applicability of feldspar IRSL₅₀ thermochronometry over a wide range of long-term (i.e. 10^5 - 10^8 years) exhumation rates. All sample locations and lithological descriptions are given in Table 1. Eight samples are distributed along the Sognefjord and its vicinity (SOG samples; western Norway) where exhumation rates are ~ 0.01 - 0.05 km Ma⁻¹ (e.g. Hendriks et al., 2007; Nielsen et al., 2009). The SOG samples were collected at varying elevations, from sea-level towards ~ 1000 -m a.s.l. low-relief surfaces (e.g. Steer et al., 2012). Fifteen samples have been collected in the Granite

Range (GRA samples; southern Alaska) where low-temperature thermochronometric data indicates long-term exhumation rates of 0.1-0.5 km Ma⁻¹ (e.g. Spotila et al., 2004; Berger et al., 2008). In the Granite Range, our sampling strategy was to collect bedrock samples across elevation profiles (Wagner and Reimer, 1972; Valla et al., 2011b). Finally, nine samples have been collected in the Yakutat terrain (YAK samples; southern Alaska), all locations at sea-level along the Disenchantment Bay - Russel Fjord. This area is tectonically active with reported long-term exhumation rates of ~1-3 km Ma⁻¹ (e.g. McAleer et al., 2009; Enkelmann et al., 2015). Our sampling strategy thus covers two orders of magnitude of long-term bedrock exhumation rates, from 0.01 km Ma⁻¹ (SOG) up to 1-3 km Ma⁻¹ (YAK). Moreover, we targeted different lithologies (sandstone, gneiss and granite/diorite samples), probing a large diversity in bedrock geochemistry and consequently also in feldspar mineralogy. Finally, we restricted our study to cold environments (mean annual surface temperatures <10 °C, Table 1) to exclude potential luminescence signal loss at surface temperature conditions in tropical to sub-tropical regions (Wu et al., 2015).

Bedrock samples were prepared under subdued orange light to ensure the extraction of feldspar minerals from the interior of the sample, where the luminescence signal is unaffected by surface bleaching. Initial sample dimensions were at least 10 cm (length/width/thickness) to allow the removal of the light-exposed sample surface (at least 2-3 cm) using a water-cooled diamond saw. The light-safe internal part was then gently hand crushed to 180-212 µm grain size using a mortar and pestle to avoid mineral grinding and potential luminescence resetting via triboluminescence (e.g. Bateman et al., 2012). All 180-212 µm grains were treated with 32% hydrochloric acid to remove potential surface carbonates and with 30% hydrogen peroxide to remove any organic component. After chemical treatment, mineral fractions were rinsed thoroughly with water to remove dust particles. Feldspar and quartz fractions were isolated by density separation using LST Fastfloat (sodium heteropolytungstate). Densities of 2.58 g cm⁻³ and 2.70 g cm⁻³ were used in order to isolate potassium-rich feldspars (K-feldspars, <2.58 g cm⁻³) from quartz and other feldspars (2.58-2.70 g cm⁻³), and from remaining heavy minerals (>2.70 g cm⁻³). For SOG samples (except SOG-22) we selected K-feldspars for analysis.

For GRA, YAK and SOG-22 samples, density separation at 2.58 g cm⁻³ yielded no light fraction, suggesting the dominance of calcium/sodium-plagioclases (Ca-/Na-feldspars), which were thus taken for further analysis. Further mineral purification was deemed unnecessary, as the contribution from quartz to IRSL₅₀ is negligible (Sohbati et al., 2013 and references

therein). Feldspar extracts were not etched before luminescence measurements (Duller, 1992), and feldspar grains were directly mounted on stainless steel discs to produce large aliquots (4 to 6 mm diameter, ~500-1500 grains per disc).

2.2. Luminescence equipment and protocols

All IRSL₅₀ measurements were made on automated Risø TL/OSL DA-20 readers (Institute of Geological Sciences, University of Bern), each possessing a ⁹⁰Sr/⁹⁰Y beta source (~0.1-0.2 Gy s⁻¹) and a heater plate (20-700 °C), with systematic instrumental uncertainties of ~1.5% for single-aliquot measurements (Bøtter-Jensen et al., 2010). Luminescence stimulation was performed with infrared (870 ±40 nm) light-emitting diodes at 90% power (delivering ~130 mW cm⁻² at the aliquot position); the emitted luminescence signal was detected through a 410-nm interference and a 2-mm Schott BG-39 filters by an EMI 9235QA photomultiplier tube. All luminescence protocols described below (adapted from Guralnik, 2015a) are fully detailed in Table 2.

Feldspar natural IRSL₅₀ signals, and their subsequent laboratory dose-response curves, were measured using the Single-Aliquot Regenerative-dose (SAR) protocol (Wallinga et al., 2000). For each sample, at least three individual aliquots were measured using a preheat at 250 °C for 60 s (to empty unstable electron traps), followed by IRSL₅₀ for 200 s (Table 2). We used laboratory beta irradiation doses from 0 up to ~2800 Gy to fully characterize the dose response curve, a test dose of 44 Gy to monitor sensitivity changes throughout the protocol, and a high-temperature optical wash (IRSL at 290 °C) between measurement cycles (Buylaert et al., 2009). Given the recent interest in and the research into “test dose plateaus” (Yi et al., in press), we believe that a test dose of 44 Gy is readily applicable to obtain equivalent doses in the ~4.4-440 Gy range (cf. Yi et al., in press), and likely even beyond. While small variations in the laboratory saturation level as a function of test dose have also been documented (Yi et al., in press), these effects are demonstrably minor and were hence considered beyond the scope of the present study.

All analysed aliquots fulfilled the standard performance criteria for recuperation (<5%), recycling ratio (<10%) and maximum test dose error (<10%). To test the efficacy of the SAR protocol parameters, a dose recovery experiment was conducted for each sample (Supplementary Table S5). To this end, two to three fresh aliquots per sample were optically bleached using IRSL₅₀ for 300 s, and were subsequently given a 200-Gy laboratory dose which

was then treated as an unknown and recovered using the same SAR protocol as the natural doses.

To explore the fading variability of $IRSL_{50}$ signals in individual bedrock samples, we conducted fading experiments on two sensitised aliquots from the dose response experiments. We used two different approaches to measure feldspar $IRSL_{50}$ fading (Table 2), namely a SAR-based long-shine method (modified from Huntley and Lamothe, 2001), and a non-SAR short-shine method (modified from Auclair et al., 2003), with the objective to test the agreement between the two protocols. In the long-shine protocol, after a given dose of 200 Gy, the luminescence signal was measured following variable delays, each lasting between 0.3 and 48 hours (Table 2). In the short-shine protocol, a dose of 72 Gy was given only once, and followed by a series of consecutive short-shine (0.1 s) measurements, separated from each other by pauses lasting between 0.1 and 55 hours. To account for signal loss due to the short optical excitation, the delayed short-shine measurements were normalized to an otherwise identical series of prompt (undelayed) measurements, done once before and once after the delayed experiment (Table 2).

2.3. Geochemistry and environmental dose rate determination

We obtained the whole-rock geochemistry (major and trace elements) of all 32 samples using inductively-coupled plasma mass spectrometry (Supplementary Tables S2-S3) to measure the dominant radionuclide content used for the environmental dose rate estimate (U, Th and K, Table 1). For a subset of 12 samples (GRA and YAK), we used the X-ray fluorescence (XRF) attachment to the Risø TL/OSL reader (Kook et al., 2012; Guralnik et al., 2015a) to determine feldspar mineralogical composition (Supplementary Table S4) and infer the internal potassium content (K_{int} , Tables 1 and S4). XRF data confirms that, in GRA and YAK samples, the aliquot mineralogy is dominated by Ca-/Na-feldspars (with up to 50% wt. quartz content), translating to an internal potassium content K_{int} between 0.2 and 4% (Supplementary Table S4). Based on these measurements, we assumed for the other GRA and YAK samples (incl. SOG-22) an averaged K_{int} of $1.9 \pm 1.4\%$ (Table 1), in agreement with previous estimates reported in the literature for Ca-/Na-feldspars (Barré and Lamothe, 2010; Sohbaty et al., 2013; Guralnik et al., 2015a). For SOG samples (except SOG-22), we used K_{int} of $12.5 \pm 1\%$ for K-feldspars following standard literature values (Huntley and Baril, 1997; Barré and Lamothe, 2010).

For each sample, the whole-rock radionuclide content (U, Th and K) and feldspar K_{int} have been converted into environmental dose rates using the conversion factors of Guérin et al. (2012), beta and gamma attenuation factors of Guérin et al. (2012) and Guérin and Mercier (2012), and alpha attenuation factors of Bell (1980) and Brennan et al. (1991). For alpha attenuation, an a -value of 0.15 ± 0.05 was used after Balescu and Lamothe (1994). The average water content in our bedrock samples was estimated at $2 \pm 2\%$ (e.g. Aitken and Xie, 1990). Cosmic dose rate was treated as negligible, being comparable in magnitude to uncertainty on the total environmental dose rate (King et al., 2016). Based on thin-section observations, and because the original grain size has been affected by rock crushing and sieving during feldspar extraction, we used a 180-2500 μm grain size range for SOG and GRA samples and a 180-1000 μm grain size range for YAK samples.

To address the two entangled uncertainties regarding (i) the grain size of the target feldspar, and (ii) the feldspar phase which is actually contributing to the IRSL_{50} signal (e.g. Sohbati et al., 2013), we considered several (2 or 4, Supplementary Table S1) dose rate scenarios following the approach of Guralnik et al. (2015a). Specifically, for all feldspars, we calculated dose rates for two end-member grain sizes (180 and 2500 μm for GRA and SOG, and 180 and 1000 μm for YAK), assuming a homogeneous distribution of internal potassium (K_{int} , Table 1) and IRSL_{50} signal contribution from the entire feldspar crystal. For Na-/Ca feldspars (GRA, YAK and SOG-22), we calculated two additional scenarios, in which all the IRSL_{50} signal originates only from K-feldspar inclusions (1 and 100 μm , $K_{int} \sim 12.5\%$) within the Na-/Ca-feldspar host crystal (e.g. Sohbati et al., 2013). We then averaged these end-member scenarios to obtain representative environmental dose rates, which ranged across more than an order of magnitude between samples, from 0.3 to 9 Gy ka^{-1} (Table 1). This range, and the associated conservative errors, reflect the variability between K-rich and K-poor feldspar phases, and the averaging across the diverse dose-rate scenarios as listed above (Supplementary Table S1).

3. Results

3.1. Dose response and fading measurements

Dose response and fading measurements are illustrated for four representative GRA, SOG and YAK samples (Figure 1), and are fully documented for all samples in the Supplementary Information (Supplementary Table S6). Analysed aliquots (4- to 6-mm diameter) usually yielded bright signals with typical IRSL_{50} decay curves. Laboratory dose-response curves (left panels in Figure 1) have been fitted using the General-Order Kinetics (GOK) growth function

(Eq. (10) in Guralnik et al., 2015c). The characteristic dose (D_0), the electron-trapping order (α), as well as the equivalent dose (D_e), are reported in Tables 3 and S5. For each aliquot, we also derived the natural IRSL₅₀ signal L_n/T_n (Figure 1) and calculated the natural trap filling level $(n/N)_{\text{nat}}$ by multiplying the apparent IRSL₅₀ intensity relative to its laboratory maximum (open stars in Figure 1) by the predicted number of occupiable traps in the laboratory (Eq. (3) in Guralnik et al., 2015a). The majority of best-fit D_0 and α values are well-constrained with typical relative uncertainties of <10% and <4%, respectively (Table 3). Natural trap filling levels $(n/N)_{\text{nat}}$ exhibit varying uncertainties (1-70%; clearly anticorrelated with $(n/N)_{\text{nat}}$ itself), and cover almost three orders of magnitude (i.e. 0.003-0.85; Table 3), thus offering a robust dataset for evaluating the net effects of environmental radiation, anomalous fading and possibly long-term cooling on the natural IRSL₅₀ intensity of bedrock feldspar.

Fading results obtained using both the long-shine (central panels in Figure 1) and short-shine (right panels in Figure 1) methods are also summarised in Table 3. Two aliquots were measured individually and averaged to calculate sample-specific g-values and standard deviations ($g_{2\text{days}}$ in %/decade; Huntley and Lamothe, 2001). Fitting the same datasets using Eq. (5) in Kars et al. (2008), we obtained an alternative measure of fading - the nearest-neighbouring hole-centre density ρ' as defined in Huntley (2006) - which we later used in the kinetic model to predict the IRSL₅₀ field saturation levels (Guralnik et al., 2015a). Inter-aliquot variability is small, confirming the good reproducibility of the fading measurement protocols, with more robust estimates for the long-shine method as explained below (Table 3). Although the short-shine protocol is less time-consuming (approximately half the time), the integrated IRSL₅₀ signal at each short-shine measurement is considerably lower, and therefore associated with larger uncertainties propagated into the resulting g-value (Table 3 and Figure 1). For feldspars with intermediate to low luminescence sensitivity (roughly one third of our dataset), the short-shine protocol yielded signals with a very low signal-to-background ratio that could not be used to derive meaningful g-values (Table 3). Figure 2 shows that, on average, fading rate estimates are in excellent agreement between the long- and the short-shine protocols. Thus, to further reduce experimental noise, in the subsequent modelling we used averaged g-values (wherever possible) from both protocols. Sample-specific IRSL₅₀ $g_{2\text{days}}$ values reveal a wide range from ~0.5 to ~45 %/decade (Table 3). K-feldspars from SOG samples exhibit $g_{2\text{days}}$ values between 3 and 11 %/decade, which is typical for K-feldspars from both bedrock and sediment origin (Huntley and Lamothe, 2001; Huntley and Lian, 2006). Na-/Ca-feldspars show

a higher variability in measured $g_{2\text{days}}$ values between samples (Figures 2 and 3), especially for GRA samples (from 0.5 to >45 %/decade).

3.2. Predictions of feldspar IRSL_{50} field saturation

To evaluate whether fading variability in our dataset is the main controlling factor for the natural IRSL_{50} intensities, we first examined the relationship of the laboratory-derived $g_{2\text{days}}$ values with the measured $(n/N)_{\text{nat}}$ values (Figure 3). Our results show a clear inverse correlation between $g_{2\text{days}}$ and $(n/N)_{\text{nat}}$ values, in agreement with previous observations (Huntley and Lian, 2006) but over a considerably wider range of fading rates. Even prior to modelling, such an inverse correlation suggests that the measured $(n/N)_{\text{nat}}$ in most of our samples is mainly controlled by anomalous fading. Only some YAK samples appear to deviate from the overall trend, with their $(n/N)_{\text{nat}}$ values markedly below field saturation (Figure 3). However, the relation in Figure 3 is only qualitative, because this empirical relationship might be influenced not only by sample-specific kinetic parameters, but also by thermal loss (which cannot be accounted for in this figure). In the following section, we test our current dataset against a recently proposed physical model predicting feldspar IRSL_{50} field saturation under anomalous fading (Guralnik et al., 2015a, incorporating the tunnelling term of Huntley, 2006, and extending the dose response of Kars et al., 2008 to non-first order).

Athermal IRSL signal loss in feldspars is caused by the quantum mechanical tunnelling of electrons from their traps towards electron holes (Huntley and Lamothe, 2001). Here we use the physical model of Huntley (2006) which is based on the assumptions that (i) the tunnelling of electrons is governed by the distance to their nearest recombination centre, and (ii) recombination centres within a feldspar crystal are randomly distributed with a given density. This model (Huntley 2006) can be used to estimate sample-specific field saturation values (e.g. Kars et al., 2008; Guralnik et al., 2015a; King et al., 2016). Following Guralnik et al. (2015a), we use a Monte-Carlo approach to numerically predict, on a sample to sample basis, the IRSL_{50} trap filling saturation values in nature $(n/N)_{\text{ss}}$ (termed “field saturation”). In the simulation (1000 runs per sample), random instances of the kinetic parameters (D_0 and α , Table 3), natural dose rates (Table 1), and ρ' (Table 3, averaged wherever possible between the long- and short-shine experiments) are drawn from their normal distributions (as given by their best-fit values). The modelling results, reported as median $(n/N)_{\text{ss}}$ and associated 68% confidence interval, are given in Table 3 and visualised against measured natural trap filling $(n/N)_{\text{nat}}$ values in Figure 4.

Figure 4 highlights the quantitative agreement between predicted $(n/N)_{ss}$ and measured $(n/N)_{nat}$ values, confirming that most of our samples are either in full field saturation, or hardly discernible from it ($<15\%$; Guralnik et al., 2015b). In other words, the natural feldspar $IRSL_{50}$ signal in the majority of our samples ($n = 25$) is controlled by electron tunnelling and does not include a detectable thermal signature of the rock exhumation history. However, more than half (66%) of YAK samples, representing the fastest-exhuming setting, exhibit $(n/N)_{nat}$ values that are considerably (more than 15%) below their predicted $(n/N)_{ss}$ values, suggesting field disequilibrium and thus high likelihood of a thermal signature that can be translated into a cooling history. The thermal signature of two additional samples (GRA-BR and SOG-17) remains in doubt, requiring further investigation that will either link them to exceptional thermal histories, or dismiss them as outliers (one in each locality).

3.3. Dose recovery results

Calculated dose recovery ratios (Supplementary Table S5) were typically within 10% of unity for all but four (GRA-13/-14/-17 and SOG-21) samples. Although imperfect, the 15-20% underestimation in the dose recovery of SOG-21 is within familiar values, even when working with a SOL2 solar simulator (e.g. Buylaert et al., 2012; King et al., 2016), and may relate to an irreversible trapping efficiency change (Kars et al., 2014). On the other hand, it is hard to blame the considerable dose recovery overestimates in GRA-13/-14/-17 (30 to 65%) on the test dose/given dose ratio (here 18%), which is in the 15-80% trust zone and certainly close to the recently recommended $\sim 30\%$ value (Yi et al., in press). To test whether the observed overestimation is due to a large residual dose, we measured the latter in GRA-13 and GRA-BR after an $IRSL_{50}$ bleach for 300 s. Low residuals of ~ 5 -8 Gy correspond to less than 2% of the equivalent doses, and to 2-3% of the recovered dose, and therefore cannot explain the observed overestimations. To test whether the problem lies in a thermal transfer (e.g. Huntley et al., 1993), we extended the dose recovery experiments further by conducting the $IRSL$ bleach at higher temperatures (100 and 200 °C). Bleaching at higher temperatures did not significantly improve dose recovery in GRA-13/-14/-17, and had no effect on a well-behaving sample (GRA-BR). Overall, unacceptable dose recovery affects only a minor subset (13%) of our dataset and might not necessarily imply that the obtained $(n/N)_{nat}$ values are not reliable (Buylaert et al., 2012). However, such dose recovery over- and underestimations do raise the question of what thermal/optical bleach is most appropriate for resetting the natural $IRSL_{50}$ intensity in feldspars that were never previously exposed to light.

4. Discussion

4.1. IRSL₅₀ and anomalous fading in bedrock feldspars

Our results confirm that feldspar IRSL₅₀ signals from bedrock extracts show consistent and reproducible luminescence characteristics using the SAR protocol. This is a clear advantage compared to OSL of bedrock quartz which, apart from a very few cases (e.g. Wu et al., 2015), is generally deemed dim and unsuitable (Jeong and Choi, 2012; Guralnik et al., 2015b). Moreover, the vast majority of our bedrock feldspars (87%, $n = 28$, Supplementary Table S5) also fulfilled the dose recovery test, which is a fundamental acceptance criterion for luminescence dating. Dose responses were successfully fitted with a General-Order Kinetics model, with a median order of 2.4 ± 0.8 , in agreement with Na-feldspars from the KTB borehole (Guralnik et al., 2015a), therefore capturing the non-linearity of the dose response curve in just a single parameter.

Fading measurements exhibit a large variability in bedrock feldspars from ~0.5 to ~45 %/decade, beyond the previously published value range (Huntley and Lamothe, 2001; Huntley and Lian, 2006). Moreover, we compared two established fading protocols in sediment dating (Huntley and Lamothe, 2001; Auclair et al., 2003) and demonstrated that both are applicable for bedrock feldspars, yielding comparable results between the two methods (Figure 2). Given the considerable experimental differences between the protocols (SAR-based and non-SAR), the agreement of the obtained results is striking. The dependence of the fading rate on the given laboratory dose (Huntley and Lian, 2006; Li and Li, 2008; Kars and Wallinga, 2009) is not evident from our dataset, although the administered laboratory doses (200 and 72 Gy for the long-shine and short-shine protocols, respectively) may be too close to detect any systematic deviation.

In our suite of bedrock feldspars, $(n/N)_{\text{nat}}$ values are mostly inversely correlated with laboratory-measured fading rates (Figure 3), qualitatively suggesting field saturation in the majority of samples. Quantitatively, we have demonstrated that field saturation can be successfully predicted across almost three orders of magnitude of IRSL₅₀ electron trap occupancy (Figure 4), using a recent extension (Guralnik et al., 2015c) of familiar electron tunnelling models (Huntley, 2006; Kars et al., 2008). Despite successful prediction of field saturation in both sedimentary feldspars (Kars et al., 2008; Kars and Wallinga, 2009) and bedrock feldspars (Guralnik et al., 2015a; King et al., 2016), the underlying model(s) and the ability to accurately measure the g-value in the laboratory remains a debateable subject (Li et

al., 2014). Here, we exploited the large variability in measured fading rates to provide a quantitative validation of our predictive model over a large range of dose rates, trap occupancies, and fading rates (Figure 4).

We also used the wide range of observed IRSL_{50} fading rates to discuss potential geochemical control, starting with the whole-rock geochemistry (Supplementary Tables S2-3) as a first-order proxy. However, we found no convincing relationship between the whole-rock trace element content (Supplementary Table S3) and feldspar IRSL_{50} fading rate. Figure 5A illustrates the lithological variability within our dataset as shown by representative major oxides (SiO_2 and Fe_2O_3), suggesting that, at first order, the fading rate would increase with decreasing SiO_2 content or increasing Fe_2O_3 (or CaO) content. Felsic lithologies (i.e. >63% wt. SiO_2 content) all show relatively low feldspar fading rates (less than 10 %/decade) regardless of the feldspar type. However, for intermediate (i.e. 52-63% wt. SiO_2 content) and mafic (i.e. <52% wt. SiO_2 content) lithologies, a large spread in fading rates is observed (Figure 5A). For a representative subset of GRA and YAK feldspar separates, we also performed XRF mineralogical determination (Supplementary Table S4). Figure 5B shows that the feldspar fading rate increases with Ca content, in agreement with previous observations (Huntley and Lian, 2006; Huntley et al., 2007). In summary, neither the whole-rock geochemistry, nor the trace elements, seem to be a good proxy for feldspar anomalous fading (although we did not apply multivariate analysis); however, our results show that the feldspar internal chemistry (i.e. Ca content, Figure 5B) may be a first-order proxy for the IRSL_{50} fading rate. Further investigations are required to better understand the potential causes of fading rate variability in bedrock feldspars; these include both internal factors (e.g. the origin of the IRSL_{50} signal in Ca-/Na-feldspars; e.g. Sohbati et al., 2013) as well as external drivers, such as the metamorphic grade or terrain age (Huntley and Lamothe, 2001), or weathering processes (Parish, 1992; Huntley, 2011).

4.2. Implications for OSL thermochronometry

Our extensive bedrock dataset offers an opportunity to discuss the applicability of feldspar IRSL_{50} in OSL thermochronometry. First, the variability of the fading rates measured for both K- and Ca-/Na-feldspars is significantly broader than those reported in previous OSL thermochronometric studies (Guralnik et al., 2015a, King et al., 2016). This reinforces the notion that anomalous fading in feldspar should be measured for each sample, as it can strongly affect the natural IRSL_{50} signal (Figure 3). Moreover, one further complexity in OSL

thermochronometry is that a sample's thermal history can be inverted only from the natural trap filling $(n/N)_{\text{nat}}$, and not from the equivalent dose alone (D_e) as in OSL sediment dating. This prevents the application of classical fading correction schemes (e.g. Huntley and Lamothe, 2001) and requires the quantitative prediction of the field saturation level of the luminescence signal (e.g. Kars et al., 2008; Li and Li, 2008) in order to screen for a potential "thermal signature" in a given sample. For most of our samples, and within a wide range of fading rates, the quantitative agreement of observed field saturation levels of the IRSL₅₀ signal with numerical results from the standard tunnelling model (Huntley, 2006), coupled with a General-Order Kinetics dose response (Guralnik et al., 2015c), confirms the need to include fading measurements in preliminary screening of samples before any thermochronometric applications using feldspar IRSL₅₀. This first-step investigation of the feldspar athermal stability would ultimately be combined with the experimental characterisation of each sample's thermal stability (Guralnik et al., 2015a; King et al., 2016), which can vary greatly between K- and Na-/Ca-feldspars (e.g. Tso et al., 1996; Li et al., 1997) and thus significantly influence thermochronometric interpretation.

Our sampling strategy, covering a wide range of exhumation rates (from 0.01 to 3 km Ma⁻¹), provides further information regarding the range of applicability of feldspar IRSL₅₀ to OSL thermochronometry in diverse natural settings. All GRA and SOG samples appear in field saturation (Figure 4), with just one outlier per locality (GRA-BR and SOG-17). SOG-17 was a translucent bedrock sample, which may immediately be suspected of suffering optical resetting (in nature and/or during sample preparation), that may have biased the observed luminescence signal to below the predicted field saturation level. GRA-BR is a spatial outlier within the GRA dataset, situated at low-elevation within a very wide and deeply-carved valley. Pending additional experimental confirmation (beyond the scope of this study), its deviation from field saturation could potentially reflect recent accelerated exhumation by efficient glacial processes. Finally, more than half of the YAK samples exhibit a thermal signature in their measured natural IRSL₅₀ signal (Table 3 and Figure 4). Samples YAK-17/-18/-50 (all with a "thermal signature") are located along a deeply-carved fjord with very high long-term exhumation rates (2-3 km Ma⁻¹; Enkelmann et al., 2015). Interestingly, the field-saturated YAK-09/-13/-15/-19 also belong to this area but are located ~20 km eastwards along another smaller fjord (Enkelmann et al., 2015). These spatial differences in IRSL₅₀ intensities would suggest spatial variations in the late-Quaternary exhumation history of this area that were not detected by higher-temperature thermochronometers (McAleer et al., 2009; Enkelmann et al., 2015) but

may be evidenced in our dataset. Finally, YAK-03/-07 exhibit exceptionally low natural trap filling levels (Figure 4), which may have been caused by hydrothermal reheating from crustal fluid circulations in a highly-fractured area (Fairweather Fault; McAleer et al., 2009), rather than by extremely rapid bedrock cooling. Further work beyond the scope of the present study will focus on constraining thermal kinetic parameters for these samples to derive their thermal histories.

In summary, our dataset confirms that feldspar IRSL₅₀ would be applicable only in very rapidly-exhuming settings ($>1 \text{ km Ma}^{-1}$; e.g. King et al., 2016) or in high-temperature environments such as boreholes or tunnels ($>35 \text{ }^{\circ}\text{C}$; Guralnik et al., 2015a). Feldspar anomalous fading, even if correctly accounted for, ultimately reduces the trapped charge capacity in a given crystal because of athermal instability, and thus initiates a faster arrival at field saturation. This might preclude the applicability of feldspar IRSL₅₀ in OSL thermochronometry. MET-pIRIR protocols (Li and Li, 2011a,b) offer the advantage of multi-thermochronometric systems in a single mineral (Qin et al., 2015), with promising outcomes in providing robust constraints on bedrock cooling rates over sub-Quaternary timescales (King et al., 2016). Other feldspar protocols such as thermally-transferred IRSL (Reimann et al., 2015) or pulsed IRSL₅₀ (although the thermal stability of pulsed IRSL₅₀ remains to be experimentally constrained; Tsukamoto et al., 2006; Roskosch et al., 2015; Jain et al., 2015b) may also be interesting to explore for thermochronometric application, providing that these protocols show good luminescence characteristics as well as low thermal stability for bedrock feldspar extracts. Alternatively, quartz protocols based on thermally-transferred OSL (Duller and Wintle, 2012), Violet Stimulated Luminescence (VSL, Ankjærgaard et al., 2013; 2015) or (red-)thermoluminescence (TL or red TL, Schmidt et al., 2015; Tang and Li, 2015) could allow for later saturation and thus a wider dating range in thermochronometry.

5. Conclusions

We report feldspar IRSL₅₀ luminescence characteristics, natural signal and anomalous fading for 32 bedrock samples collected from a large variety of lithologies and exhumation settings. Our results show that feldspar IRSL₅₀ signals from bedrock extracts are bright and reproducible using established protocols in OSL dating. We measured a large variability in IRSL₅₀ fading rates between samples. Our results show that it is essential to measure anomalous fading and account for it on a sample-to-sample basis for OSL thermochronometry, as it may otherwise obscure any thermochronometric information. Furthermore, we exploited

the wide range of observed fading rates to demonstrate the ability of electron tunnelling models in quantitatively predicting the feldspar IRSL₅₀ level in field saturation for rather diverse environmental conditions. While the potential causes for feldspar IRSL₅₀ fading still require further investigation, our observations point towards a significant influence of Ca content in feldspar on its anomalous fading rate. Finally, our results suggest that the applicability of feldspar IRSL₅₀ in OSL thermochronometry would be limited to rapidly-exhuming settings (i.e. >1 km Ma⁻¹) or warm subsurface environments (i.e. >35 °C).

Acknowledgments

We thank S. Tsukamoto and S.-H. Li for constructive and helpful reviews, A.S. Murray and M. Jain (Nordic Laboratory for Luminescence dating / DTU Risø) for hosting an early pilot study in 2010, M. Kook for XRF measurements, and G. King and D. Sanderson for fruitful discussions and feedback. P.G.V. acknowledges funding from the SNSF (grant Ambizione PZ00P2_148191) and the ETH-Zürich COFUND. S.E.L. receives funding from the National Cooperative for the Disposal of Radioactive Waste (Nagra). F.H. and B.G. were funded through SNSF grants PP00P2-38956 and 200021-127127. B.G. acknowledges support from the NWO VENI grant 863.15.026. P.S. acknowledges the financial and scientific support from R.S. Huismans, the Statoil Earth System Modelling project and the University of Bergen.

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Appendix A. Supplementary material

Supplementary material, including details about the dose rate scenarios (Table S1), full geochemical reports (Tables S2-S4), supplementary experimental results (Table S5) and all sample raw data (Table S6) related to this article can be found on-line at ...

Figure captions

Figure 1. IRSL₅₀ dose-response curves (left panels) and fading experiments (central and right panels) for representative samples GRA-08 (A-B-C), GRA-12 (D-E-F), SOG-06 (G-H-I) and YAK-09 (J-K-L). See Tables 2 and 3 for experimental details and results, respectively. The dose-response curves (left panels) have been obtained following the single aliquot regenerative-dose (SAR) protocol for 3 individual aliquots. The natural luminescence signal (white stars) and each regenerative dose (black circles) are normalised to a given test dose, and used to build the dose-response curve (fitted by the GOK approach, black line and grey envelope; α : electron-trapping order, D_0 : characteristic dose) and to calculate the sample natural luminescence level (L_n/T_n) and natural trap filling $(n/N)_{nat}$ (Guralnik et al., 2015a). Fading experiments, following the long-shine (central panels) and short-shine (right panels) protocols, were used to derive the sample fading rate (sample-specific g_{2days} are given by the slope of the black lines normalized to 2 days; Huntley and Lamothe, 2001).

Figure 2. Comparison of fading rate measurements (g_{2days}) using the short-shine and long-shine methods (Table 3). The symbols represent individual GRA (black circles), YAK (grey squares) and SOG (open triangles) samples. The dashed line defines the 1:1 relationship. The inset shows a zoom-in to the low-fading rate region (dashed box in main panel).

Figure 3. Natural trap filling $(n/N)_{nat}$ and g_{2days} values (long-shine method, Table 3) for GRA (black circles), YAK (grey squares) and SOG (open triangles) samples. Samples in field saturation define a well-characterised inverse relationship between laboratory-measured fading rates (g_{2days}) and the natural trap filling $(n/N)_{nat}$ values over a wide range of fading rates (from ~0.5 to ~45 %/decade).

Figure 4. Predicted field saturation $(n/N)_{ss}$ vs. measured natural trap filling $(n/N)_{nat}$ for GRA (circles), YAK (squares) and SOG (triangles) samples (A: lineal plot; B: logarithmic plot). Symbols are coloured according to the individual sample g_{2days} values (see legend). Most of

analysed samples plot along the 1:1 relationship (dashed line), validating the approach of Guralnik et al. (2015a) to numerically predict sample field saturation (after Huntley, 2006 and Kars et al., 2008). Some samples have $(n/N)_{\text{nat}}$ below $0.85(n/N)_{\text{ss}}$ (dotted line), bearing potential thermochronometric information. See text for details and discussion.

Figure 5. Feldspar geochemistry and laboratory-measured fading rate. (A) Whole-rock representative oxides showing the lithological variability between analysed samples (Supplementary Table S2). (B) Feldspar mineralogical compositions for a subset of samples (Ab: albite; An: anorthite; Or: orthoclase; Supplementary Table S4). GRA: circles; YAK: squares; SOG: triangles. Symbols are coloured according to the respective sample's $g_{2\text{days}}$ value (see legend). See text for discussion.

Table 1. Overview of sample locations, lithology and climate, alongside measured radionuclides and calculated natural dose rate.

Sample	Latitude/ Longitude (°N)/(°W)	Elevation (m)	Lithology	Mean Annual Temperature ¹ (°C)	Whole-rock geochemistry ²			Feldspar minimum K_{int} ³ (%)	Natural dose rate D_{rate} ⁴ (Gy ka ⁻¹)
					U (ppm)	Th (ppm)	K (wt. %)		
<i>Granite Range (southern Alaska)</i>									
GRA-03	60.853/ 141.451	1419	Diorite	-5.8	0.4	1.1	0.79	0.7 ^a	1.12 (±0.24)
GRA-04	60.832/141.467	1761	Diorite	-4.5	0.2	0.6	0.82	1.7 ^a	1.22 (±0.23)
GRA-05	60.835/141.469	1995	Diorite	-6.3	0.4	1.1	0.86	1.4 ^a	1.31 (±0.14)
GRA-06	60.841/141.501	2388	Diorite	-6.8	0.3	0.5	0.76	1.0 ^a	1.06 (±0.14)
GRA-08	60.905/141.970	2188	Diorite	-7.3	0.3	0.8	0.97	1.4 ^a	1.34 (±0.13)
GRA-09	60.901/141.954	1691	Diorite	-5.7	0.1	0.1	0.50	0.8 ^a	0.70 (±0.15)
GRA-10	60.900/141.934	1412	Diorite	-4.2	0.3	1.1	2.02	3.9 ^a	2.72 (±0.53)
GRA-11	60.908/141.919	970	Diorite	-3	0.6	1.1	0.81	1.4 ^a	1.33 (±0.16)
GRA-12	60.991/142.763	1554	Schist	-2.5	0.3	0.5	1.03	1.9 ^b	1.46 (±0.22)
GRA-13	61.003/142.741	1253	Granitoïde	-1.5	1.2	9.4	0.72	1.9 ^b	2.37 (±0.61)
GRA-14	60.991/142.782	1979	Granitoïde	-4.2	0.9	1.3	1.09	1.9 ^b	1.79 (±0.20)
GRA-17	60.943/142.475	1818	Granitoïde	-6.1	0.3	0.5	0.61	1.9 ^b	1.11 (±0.35)
GRA-18	60.957/142.378	1395	Schist	-4.1	0.1	0.3	0.61	1.9 ^b	1.02 (±0.39)
GRA-19	61.008/142.362	720	Granitoïde	-1.4	1.6	3.6	0.73	1.9 ^b	1.95 (±0.39)
GRA-BR	61.109/142.420	392	Monzogranite	-0.2	1.6	3.9	2.17	4.2 ^a	3.62 (±0.49)
<i>Sognefjord (western Norway)</i>									
SOG-02	61.096/-5.681	33	Granitic gneiss	7.4	1.8	10.4	4.18	12.5 ^c	8.45 (±3.44)
SOG-06	61.155/-6.461	883	Granitic gneiss	2.4	2.2	8.3	5.13	12.5 ^c	9.01 (±3.01)
SOG-10	61.146/-6.542	47	Granitic gneiss	6.5	1.1	10.3	4.76	12.5 ^c	8.68 (±3.24)
SOG-17	61.233/-7.593	961	Granodiorite	1.1	0.2	0.2	4.13	12.5 ^c	7.51 (±3.94)
SOG-21	61.239/-7.664	10	Anorthosite	3.5	0.7	10.5	5.15	12.5 ^c	8.87 (±3.10)
SOG-22	61.263/-7.867	1274	Gabbro	-0.9	0.1	0.8	0.52	1.9 ^b	1.00 (±0.40)
SOG-25	61.272/-7.765	48	Anorthosite	4.8	0.6	5.5	5.30	12.5 ^c	8.65 (±3.19)
SOG-38	61.151/-6.728	21	Granitic gneiss	7.0	0.4	3.7	5.13	12.5 ^c	8.40 (±3.34)

<i>Yakutat terrane (southern Alaska)</i>									
YAK-03	59.813/138.912	5	Diorite	0.3	0.1	0.3	0.18	0.2 ^a	0.34 (±0.17)
YAK-07	59.843/139.162	2	Diorite	1.6	0.1	0.1	0.04	1.9 ^b	0.33 (±0.22)
YAK-09	59.825/138.986	0	Diorite	2.3	2.5	4.4	1.72	2.2 ^a	3.24 (±0.70)
YAK-13	59.647/139.296	2	Sandstone	3.5	1.7	6.3	1.37	1.9 ^b	2.80 (±0.63)
YAK-15	59.851/139.313	1	Sandstone	3.0	2	4.8	1.53	1.9 ^b	2.90 (±0.63)
YAK-17	59.974/139.450	2	Sandstone	2.2	1.4	3.9	0.62	1.9 ^b	1.74 (±0.39)
YAK-18	59.967/139.531	1	Sandstone	2.9	1.2	3.5	1.87	1.9 ^b	2.80 (±0.44)
YAK-19	59.597/139.342	2	Sandstone	3.2	1.6	3.2	1.89	4.2 ^a	3.12 (±0.37)
YAK-50	59.946/139.620	0	Sandstone	2.6	1.2	3.5	1.11	1.9 ^b	2.08 (±0.37)

Notes. ¹ Derived from Worldclim database (Hijmans et al., 2005). ² Measured on a Perkin Elmer Sciex ELAN 6100/9000 ICP-MS, with an estimated analytical precision of 2.3 (U), 3.4% (Th) and 1.2% (K). See Supplementary Tables S2-S3 for all elements and further details. ³ Measured or estimated as follows: ^a measured using the Risø XRF attachment (Kook et al., 2012); ^b assumed value for Na/Ca-feldspars (average of the Risø XRF data above) ^c assumed value for K-feldspars (Huntley and Baril, 1997; Barré and Lamothe, 2010) ⁴ For calculation of feldspar dose rates, we used conversion factors of Guérin et al. (2012), water content of 2(±2)%, and negligible cosmic radiation. The reported dose rates cover two (or four) end-member scenarios after Guralnik et al. 2015a (see Table S1). For K-feldspars, two end-member grain sizes were considered (180 and 1000 or 2500 µm); for Na/Ca-feldspars, two additional scenarios, in which the entire luminescence signal originates from K-feldspar inclusions (1 or 100 µm) were further included (Table S1).

Table 2. Laboratory protocols for luminescence measurements of feldspar IRSL₅₀.

Dose response SAR protocol ^a	Anomalous fading	
	SAR protocol, (Long-shine method ^b)	Short-shine method ^c
1. Give dose D_x	1. IRSL at 290 °C (200 s)	1. IRSL at 290 °C (200 s)
2. Preheat (250 °C for 60 s)	2. Give dose D_{fading}	2. Give dose D_{fading}
3. IRSL at 50 °C (200 s) $\rightarrow L_x$	3. Preheat (250 °C for 60 s)	3. Preheat (250 °C for 60 s)
	4. Hold at room T °C for t_x s	4. Hold at room T °C for t_x s
	5. IRSL at 50 °C (200 s) $\rightarrow L_x$	5. IRSL at 50 °C (0.1 s) $\rightarrow L_x$
		6. Return to step 4
4. Give test dose D_{test}	6. Give test dose D_{test}	7. IRSL at 290 °C (200 s)
5. Preheat (250 °C for 60 s)	7. Preheat (250 °C for 60 s)	8. Give test dose D_{test}
6. IRSL at 50 °C (200 s) $\rightarrow T_x$	8. IRSL at 50 °C (200 s) $\rightarrow T_x$	9. Preheat (250 °C for 60 s)
7. IRSL at 290 °C (200 s)		10. IRSL at 50 °C (0.1 s) $\rightarrow T_x$
8. Return to step 1	9. Return to step 1	11. Return to step 10

Protocols adapted from Guralnik et al. (2015a). ^a The irradiation doses of the SAR dose response protocol (Wallinga et al., 2000) were $D_x = [0, 22, 44, 87, 175, 349, 698, 1396, 2792, 0, 87]$ Gy, and the test dose was $D_{test} = 44$ Gy. The first and the penultimate zero-doses (no irradiation) extract the natural dose and signal recuperation, respectively. ^b The irradiation and test doses of the long-shine fading experiment (Huntley and Lamothe, 2001) were $D_{fading} = 240$ Gy and $D_{test} = 44$ Gy, respectively, and the holding times were $t_x \approx [1200, 1200, 1200, 136660, 1200, 7970, 1200, 143450, 1200, 1200, 1200]$ s. ^c The short-shine fading experiment (Auclair et al., 2003) consisted of a single irradiation dose $D_{fading} = 72$ Gy which was given, followed by delay times of $t_x \approx 10^{2.34+0.37x}$ s, where $x = [1, 2, \dots, 8]$. The test dose was $D_{test} = 72$ Gy. L_x and T_x in both SAR protocols (left and central column) were derived by integrating the first 15 s of the stimulation curve (signal) after subtraction of the last 100 s (background). L_x and T_x of the short-shine experiment (right column) were derived by integrating the entire stimulation period of 0.1 s, after subtraction of a 20 ms background (immediately following the IRSL stimulation). Note: All preheats were done with a linear heating rate of 5 °C s⁻¹.

Table 3. Best-fit parameters and model results for feldspar IRSL₅₀.

Sample	Fitted kinetic parameters						Interpretation	
	D ₀ (Gy)	Electron-trapping order (α)	Long-shine g _{2days} (%/decade)	Long-shine log ₁₀ p'	Short-shine g _{2days} (%/decade)	Short-shine log ₁₀ p'	Observed (n/N) _{nat}	Predicted (n/N) _{ss}
GRA-03	1208 (±59)	4	41.3 (±1.2)	-4.69 (±0.01)	47.7 (±4.5)	-4.68 (±0.03)	0.003 (±0.003)	0.002 (±0.0002)
GRA-04	1208 (±75)	4	28.5 (±0.7)	-4.81 (±0.01)	Low signal-to-noise ratio		0.007 (±0.005)	0.007 (±0.001)
GRA-05	550 (±71)	2.56 (±0.82)	28.0 (±1.0)	-4.82 (±0.01)	23.5 (±3.1)	-4.89 (±0.04)	0.017 (±0.003)	0.017 (±0.004)
GRA-06	612 (±29)	1.64 (±0.26)	36.2 (±1.4)	-4.74 (±0.01)	23.3 (±3.2)	-4.89 (±0.05)	0.008 (±0.002)	0.012 (±0.002)
GRA-08	599 (±31)	1.60 (±0.28)	33.8 (±0.8)	-4.75 (±0.01)	29.9 (±6.3)	-4.82 (±0.07)	0.015 (±0.003)	0.009 (+0.002/-0.003)
GRA-09	718 (±115)	2.11 (±0.86)	23.6 (±0.6)	-4.88 (±0.01)	22.2 (±3.2)	-4.92 (±0.05)	0.016 (±0.007)	0.022 (±0.006)
GRA-10	519 (±36)	1.67 (±0.38)	12.9 (±0.3)	-5.10 (±0.01)	14.3 (±1.7)	-5.07 (±0.04)	0.120 (±0.027)	0.098 (+0.016/-0.012)
GRA-11	606 (±73)	4	11.1 (±0.6)	-5.16 (±0.02)	Low signal-to-noise ratio		0.074 (±0.015)	0.113 (+0.011/-0.012)
GRA-12	545 (±10)	2.92 (±0.12)	13.9 (±0.3)	-5.08 (±0.01)	14.2 (±1.3)	-5.07 (±0.03)	0.092 (±0.005)	0.074 (+0.009/-0.007)
GRA-13	132 (±3)	1.62 (±0.08)	0.5 (±0.3)	-6.50 (±0.26)	Low signal-to-noise ratio		0.845 (±0.024)	0.903 (+0.054/-0.023)
GRA-14	243 (±4)	1.76 (±0.07)	3.4 (±0.3)	-5.65 (±0.04)	3.7 (±1.1)	-5.60 (±0.13)	0.589 (±0.017)	0.512 (+0.078/-0.061)
GRA-17	171 (±2)	3.02 (±0.11)	3.0 (±0.6)	-5.70 (±0.09)	Low signal-to-noise ratio		0.432 (±0.015)	0.567 (+0.041/-0.093)
GRA-18	245 (±3)	1.99 (±0.07)	7.5 (±0.4)	-5.32 (±0.02)	Low signal-to-noise ratio		0.276 (±0.013)	0.260 (+0.013/-0.020)
GRA-19	875 (±36)	4	9.1 (±2.0)	-5.24 (±0.09)	Low signal-to-noise ratio		0.097 (±0.014)	0.153 (+0.069/-0.037)
GRA-BR	433 (±11)	4	5.0 (±0.5)	-5.49 (±0.04)	Low signal-to-noise ratio		0.234 (±0.036)	0.376 (+0.028/-0.037)
SOG-02	688 (±16)	1.57 (±0.13)	8.6 (±0.2)	-5.26 (±0.01)	Low signal-to-noise ratio		0.172 (±0.010)	0.229 (+0.015/-0.017)
SOG-06	613 (±8)	1.41 (±0.07)	4.3 (±0.1)	-5.54 (±0.01)	5.1 (±1.0)	-5.47 (±0.08)	0.409 (±0.044)	0.444 (+0.09/-0.019)
SOG-10	588 (±10)	1.75 (±0.09)	4.6 (±0.1)	-5.52 (±0.01)	4.6 (±0.7)	-5.52 (±0.06)	0.403 (±0.050)	0.444 (+0.010/-0.013)
SOG-17	384 (±6)	1.87 (±0.09)	3.7 (±0.1)	-5.61 (±0.01)	5.2 (±0.3)	-5.47 (±0.02)	0.294 (±0.019)	0.448 (+0.006/-0.005)
SOG-21	717 (±10)	2.21 (±0.08)	8.5 (±0.1)	-5.27 (±0.01)	11.1 (±0.5)	-5.16 (±0.02)	0.164 (±0.009)	0.180 (+0.011/-0.010)
SOG-22	236 (±5)	1.30 (±0.06)	4.0 (±0.3)	-5.57 (±0.03)	3.2 (±0.4)	-5.67 (±0.05)	0.519 (±0.022)	0.480 (+0.049/-0.023)
SOG-25	567 (±12)	2.11 (±0.13)	5.2 (±1.3)	-5.47 (±0.10)	4.8 (±2.1)	-5.50 (±0.18)	0.359 (±0.016)	0.408 (+0.057/-0.108)
SOG-38	552 (±9)	2.19 (±0.11)	3.9 (±1.1)	-5.58 (±0.12)	4.5 (±0.5)	-5.52 (±0.04)	0.439 (±0.020)	0.446 (+0.049/-0.027)

YAK-03	349 (± 15)	1.72 (± 0.23)	12.6 (± 1.6)	-5.12 (± 0.05)	Low signal-to-noise ratio		0.034 (± 0.008)	0.094 (+0.039/-0.020)
YAK-07	276 (± 6)	1.51 (± 0.08)	4.1 (± 0.4)	-5.56 (± 0.04)	4.1 (± 2.3)	-5.56 (± 0.24)	0.011 (± 0.008)	0.443 (+0.112/-0.132)
YAK-09	194 (± 3)	1.70 (± 0.06)	3.3 (± 0.3)	-5.65 (± 0.04)	3.1 (± 0.3)	-5.67 (± 0.05)	0.549 (± 0.018)	0.577 (+0.022/-0.039)
YAK-13	431 (± 13)	2.44 (± 0.21)	5.5 (± 0.4)	-5.44 (± 0.03)	7.4 (± 1.7)	-5.32 (± 0.09)	0.250 (± 0.023)	0.289 (+0.046/-0.025)
YAK-15	271 (± 11)	2.32 (± 0.29)	4.1 (± 0.5)	-5.56 (± 0.05)	Low signal-to-noise ratio		0.427 (± 0.071)	0.447 (+0.045/-0.024)
YAK-17	443 (± 22)	4	6.8 (± 0.5)	-5.36 (± 0.03)	Low signal-to-noise ratio		0.125 (± 0.036)	0.254 (+0.017/-0.020)
YAK-18	405 (± 7)	2.24 (± 0.11)	8.1 (± 0.4)	-5.29 (± 0.02)	5.5 (± 1.0)	-5.44 (± 0.08)	0.204 (± 0.007)	0.282 (+0.024/-0.020)
YAK-19	905 (± 13)	4	7.2 (± 0.6)	-5.33 (± 0.03)	Low signal-to-noise ratio		0.121 (± 0.014)	0.233 (+0.019/-0.022)
YAK-50	390 (± 7)	2.36 (± 0.12)	4.1 (± 0.3)	-5.57 (± 0.03)	3.3 (± 1.1)	-5.66 (± 0.15)	0.124 (± 0.005)	0.476 (+0.090/-0.038)

Notes. Dose response curves (with recycling ratio within 10% of unity) have been fitted using the General-Order Kinetics GOK function (Guralnik et al., 2015a,c), to obtain the characteristic dose D_0 and the kinetic order α (note: a limiting value of $\alpha = 4$ was adopted for a few datasets which did not exhibit sufficient flattening off of the luminescence signal at high doses). Fading decay curves were fitted (a) following Huntley and Lamothe (2001) to obtain the g -value (normalized to 2 days); (b) following Huntley (2006) to obtain the density of recombination centres ρ' (reported as $\log_{10}\rho'$), assuming an escape frequency factor of $3 \times 10^{15} \text{ s}^{-1}$ (Huntley, 2006). Results of the short-shine method were accepted only when there existed two replicates with acceptable signal to noise ratio (luminescence signal ten times higher than background). Natural trap filling (n/N_{nat} ; reported as mean and std. dev.) and predicted field saturation (n/N_{ss} ; reported as median and 68% conf. int.) were derived following Guralnik et al. (2015a), the latter using a Monte-Carlo approach (1000 iterations, each with randomized input of the kinetic parameters and the dose rate; note that wherever possible, an average ρ' between the long- and short-shine experiments was used). Laboratory dose rates vary between 0.099 and 0.207 Gy s^{-1} for the different instruments used.

